

⁵⁷Fe Isomer Shifts and the Problem of Calibration

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Z. Naturforsch. **42a**, 211–212 (1987);
received November 24, 1986

The ⁵⁷Fe isomer shift calibration problem is discussed, and critical comments concerning the paper “⁵⁷Fe Isomer Shift Calibration Experiment” by Daniel et al. (1985) are given. We consider $-\langle \Delta r^2 \rangle = (20 \pm 3) \cdot 10^{-3} \text{ fm}^2$ at the moment as the most reliable estimate for the change of nuclear radius during gamma absorption; this value has been derived by measuring changes in isomer shift δ and calculating corresponding changes in electron contact density $\varrho(0)$ according to $\Delta \delta = C \langle \Delta r^2 \rangle \Delta \varrho(0)$.

A reliable calibration of the isomer shift δ is the prerequisite for deriving contact densities $\varrho(0)$ from Mössbauer spectra. It is therefore important to design experiments with the aim to measure directly the change of nuclear radius $\langle \Delta r^2 \rangle$ during gamma resonance absorption. Such experiments have been performed for ⁵⁷Fe on the basis of life time variations in the electron capture decay of ⁵²Fe [1] or on the basis of conversion electron spectroscopy with ⁵⁷Co(⁵⁷Fe) sources [2]. However, the results, which have been obtained with these two methods, $-\langle \Delta r^2 \rangle = (33 \pm 3) \cdot 10^{-3} \text{ fm}^2$ [1] and $-\langle \Delta r^2 \rangle < 9 \cdot 10^{-3} \text{ fm}^2$ [2], respectively, do not agree at all. Other authors, including ourselves, strived for this goal by measuring changes in δ and calculating corresponding changes in $\varrho(0)$ [3, 4]:

$$\Delta \delta = C \langle \Delta r^2 \rangle \Delta \varrho(0). \quad (1)$$

If δ is given in mm s^{-1} , $\langle \Delta r^2 \rangle$ in 10^{-3} fm^2 , and $\varrho(0)$ in a_0^{-3} , the constant C takes the value 0.011 for ⁵⁷Fe. From the combined experimental ($\Delta \delta$) and calculational ($\Delta \varrho(0)$) procedure we have derived $\langle \Delta r^2 \rangle$ for a large variety of iron-containing compounds, i.e. $-\langle \Delta r^2 \rangle = (20 \pm 3) \cdot 10^{-3} \text{ fm}^2$ [4–6].

Due to its importance it is desirable to have a widely accepted isomer shift calibration. In our view there are serious drawbacks which cause us to

question the value $-\langle \Delta r^2 \rangle < 9 \cdot 10^{-3} \text{ fm}^2$ (at 80% confidence) [2], due to the following reasons:

(i) The absolute value of the Fe 4s contact density $\varrho_{4s}(0)$ was obtained in [2] from the experimental ratio α_{4s}/α_{3s} and the calculated relativistic Fe 3s value ($\varrho_{3s}(0) \sim 180 a_0^{-3}$), i.e.

$$\text{from } \alpha_{4s}/\alpha_{3s} = 0.061 \pm 0.005 \text{ the value} \\ \varrho_{4s}(0) = (10.98 \pm 0.90) a_0^{-3} \text{ for } ^{57}\text{Fe}/\text{Au}, \quad (2a)$$

and

$$\text{from } \alpha_{4s}/\alpha_{3s} = 0.17 \pm 0.10 \text{ the value} \quad (2b) \\ \varrho_{4s}(0) = (30.6 \pm 18.0) a_0^{-3} \text{ for } ^{57}\text{Fe}/\text{graphite}.$$

With the experimental isomer shifts 0.65 mm s^{-1} for ⁵⁷Fe/Au and 0.25 mm s^{-1} for ⁵⁷Fe/graphite and the values (2a, b) we derive from (1)

$$-\langle \Delta r^2 \rangle = 1.85 \cdot 10^{-3} \text{ fm}^2.$$

Comparison with the upper limit of $9 \cdot 10^{-3} \text{ fm}^2$ given in [2] shows that the error margin is about $7 \cdot 10^{-3} \text{ fm}^2$, which means that the confidence grows only slowly when one raises the upper limit. By exhausting the full uncertainty of (2a, b) one arrives at

$$-\Delta \langle r^2 \rangle \lesssim 50 \cdot 10^{-3} \text{ fm}^2,$$

without even having considered any uncertainties of isomer shifts, because they were not reported in [2]!

(ii) The values for $\varrho_{4s}(0)$ in (2a) and (2b) are larger than the corresponding value $\varrho_{4s}(0) = 8.36 a_0^{-3}$ for Fe⁰ ($3d^6 4s^2$) as derived from Dirac-Fock calculations [4], the same calculations which yield $\varrho_{3s}(0) \sim 180 a_0^{-3}$ from above. This implies that the isomer shifts of ⁵⁷Fe/Au and ⁵⁷Fe/graphite should be even more negative than that of Fe⁰ in solid noble gas ($\delta_{x-\text{Fe}} \sim -0.75 \text{ mm s}^{-1}$ [7]), which is by far not the case.

(iii) The value for $\langle \Delta r^2 \rangle$ was estimated by Daniel et al. [2] on the basis of (1) by assuming $\Delta \varrho(0) = \Delta \varrho_{4s}(0)$. But this assumption is not necessarily satisfied, as shown in a conversion electron study of $\varrho_{4s}(0)$ of Fe-impurity atoms in transition and noble metals [8]. In fact, we have found from electronic structure calculations [4, 5] that core contributions to $\Delta \varrho(0)$ can be significantly larger than $1 a_0^{-3}$.

In conclusion, we consider $-\langle \Delta r^2 \rangle = (20 \pm 3) \cdot 10^{-3} \text{ fm}^2$ at the moment as the most reliable estimate, because we do not regard the result

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reported in [2] as conclusive, and because the value $-\langle \Delta r^2 \rangle = (33 \pm 3) \cdot 10^{-3} \text{ fm}^2$ given by Meykens et al. [1] seems to be overestimated by about 30 per cent due to reasons described in [2].

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